

# High Ionic Conductivity of $\text{Li}_2\text{O}:\text{B}_2\text{O}_3:\text{V}_2\text{O}_5$

Kanchan Gaur and H. B. Lal

Department of Physics, University of Gorakhpur,  
Gorakhpur-273009 (India)

Z. Naturforsch. **47a**, 1011–1012 (1992);  
received March 5, 1992

Lithium ion conducting solids have been of great interest in recent years [1–5]. We have also been active in this direction [6–8]. This paper reports on the electrical transport and phase transition of  $\text{Li}_2\text{O}:\text{B}_2\text{O}_3:\text{V}_2\text{O}_5$ . The compound was prepared by solid state reaction technique. The measurements of the electrical conductivity ( $\sigma$ ), thermoelectric power ( $S$ ), dielectric constant ( $K$ ) and molar magnetic susceptibility ( $\chi_M$ ) are performed from 500 to 1110 K. The details of these measurements are described in [9, 10]. The  $\sigma$  and  $S$  variations are shown in Fig. 1 as  $\log \sigma T$  and  $S$  vs.  $1/T$  plots. Both the plots show three linear regions, namely, 500–850 K, 900–960 K, and 980–1110 K. The activation energy ( $E_a$ ) and heat of transport ( $Q$ ) for the third temperature range are 0.39 eV and 0.58 eV, respectively.  $\sigma$  jumps by a factor of 180 around 855 K, and again jumps by a factor of 4 around 970 K. Thus this material seems to have two phase transition temperatures: one around 855 K and the other at 970 K. The sign of  $S$  indicates that current is carried by positive charge carriers. The time variation of the dc electrical conductivity (the results are not shown) shows that  $\sigma$  is completely ionic in the temperature range 980–1110 K. Ions are also the dominant charge carriers at lower temperatures but in general  $\sigma$  is mixed (ionic and electronic). Lithium seems to be the mobile ion of the solid.  $\sigma$  in the temperature range 980–1110 K is high, being  $150 \Omega^{-1} \text{m}^{-1}$ , and it is totally ionic, hence this is the superionic phase of the solid. The phase transitions around 855 K and 970 K are also reflected in  $K$  and  $\chi_M$  vs.  $T$  plots, as shown in Figs. 2 and 3, respectively. The room temperature  $K$  value is not high, being about 30 around 300 K. It increases slowly up to 700 K but rises steeply above 855 K and becomes of the order of  $3.2 \times 10^7$  at 900 K. It drops by a factor of 3 around 970 K and then remains almost constant. It

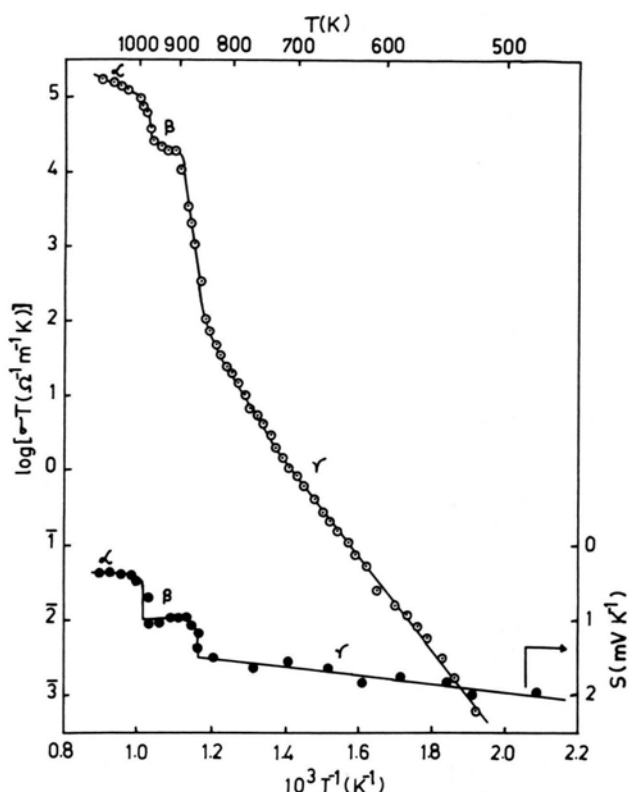


Fig. 1. Plots of  $\log \sigma T$  and thermoelectric power ( $S$ ) against inverse temperature ( $T^{-1}$ ).

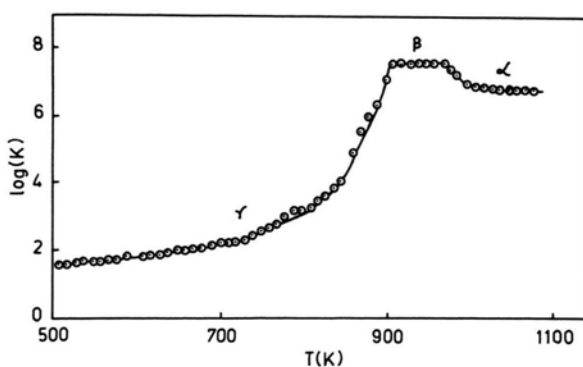


Fig. 2. Plot of logarithm of dielectric constant ( $\log K$ ) against temperature ( $T$ ).

is seen from the  $\chi_M$  vs.  $T$  plot that  $\chi_M$  is negative and is almost independent of temperature up to 855 K but shows a downward drop above this temperature. Around 970 K,  $\chi_M$  jumps abruptly and then increases slowly with temperature. The reason for the negative

Reprint requests to Dr. H. B. Lal, Department of Physics,  
University of Gorakhpur, Gorakhpur 273009/India

0932-0784 / 92 / 0900-1011 \$ 01.30/0. – Please order a reprint rather than making your own copy.



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland Lizenz.

Zum 01.01.2015 ist eine Anpassung der Lizenzbedingungen (Entfall der Creative Commons Lizenzbedingung „Keine Bearbeitung“) beabsichtigt, um eine Nachnutzung auch im Rahmen zukünftiger wissenschaftlicher Nutzungsformen zu ermöglichen.

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

On 01.01.2015 it is planned to change the License Conditions (the removal of the Creative Commons License condition "no derivative works"). This is to allow reuse in the area of future scientific usage.

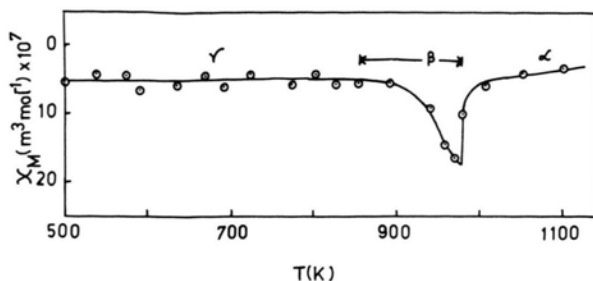


Fig. 3. Plot of molar magnetic susceptibility ( $\chi_M$ ) against temperature ( $T$ ).

### Acknowledgement

Authors are thankful to CSIR, India for financial assistance.

- [1] S. Pizzini, *J. Appl. Electrochem.* **1**, 153 (1971).
- [2] M. Lazzari and B. Scrosati, *J. Power Sources* **1**, 333 (1976/1977).
- [3] R. A. Huggins, *Electrochimica Acta* **22**, 773 (1977).
- [4] B. B. Owens, *Int. Symp. Solid Ionic and Ionic Electronic Conductors*, Rome, Sept. (1976).
- [5] H. V. Venkatesetty, *Lithium Battery Technology*, John Wiley and Sons, Inc. (1984).
- [6] Kanchan Gaur, A. J. Phatak, and H. B. Lal, *J. Mat. Sc.* **23**, 4257 (1988).
- [7] H. B. Lal, Kanchan Gaur, and A. J. Phatak, *J. Phys. D.* **22**, 305 (1989).
- [8] H. B. Lal, Kanchan Gaur, and A. J. Phatak, *J. Mat. Sc.* **24**, 1159 (1989).
- [9] Kanchan Gaur, Ph.D. Thesis, Gorakhpur University, Gorakhpur, India (1984).
- [10] A. K. Tripathi, Ph.D. Thesis, Gorakhpur University, Gorakhpur, India (1981).

value, drop and rise in  $\chi_M$  is that this material does not contain magnetic ions and  $\chi_M$  for such solids depends upon their bonding configuration. The onset of disordering at the phase transition disrupts the bonding configuration of the solid. This leads to specific changes in  $\chi_M$  at this temperature.  $Q$  is larger than  $E_a$  in the superionic phase. This is against the trend  $E_a > Q$  observed for  $\text{Ag}^+$  and  $\text{Cu}^+$  active superionic solids. This indicates that the mechanism of ion conduction in this solid differs from that usual ionic solids.